

# Contributions to Forest Fire Simulations: Mathematical Models, Numerical Methods and GIS Integration

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April 16, 2008



## 1 INTRODUCTION

- Propagation models
- Combustion models

## 2 SIMPLIFIED PHYSICAL MODEL

- The goal
- A Simple Physical Model
- Enthalpy operator

## 3 NUMERICAL METHOD

- Time Integration
- Solution at each time step

## 4 Non Local Radiation Model

- Characteristic Method

## 5 Wind Model

## 6 Simulations

- wind model example
- Physical and numerical data



# Clasification

## Models

### PROPAGATION

Position of the fire front

### COMBUSTION

Modelization of physics

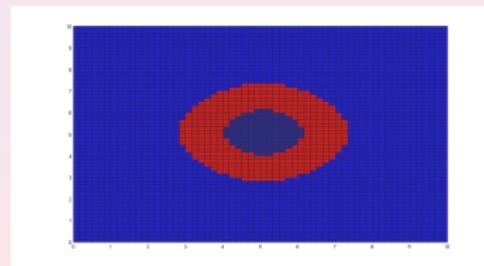
models with increasing complexity:

- Cellular Automata
- Geometric
- Empiric models
- Reaction Diffusion Convection



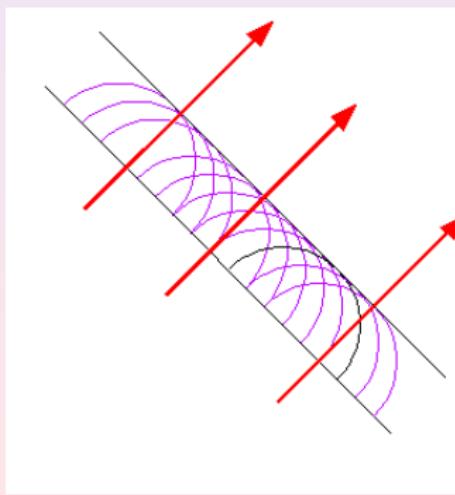
# Cellular Automata

- Cell with several states (burning, burnt, not burning)
- Transition probabilities as a function of the neighbours cells
  - Advantages:  $\Rightarrow$  Fast computation
  - Disadvantages:
    - It modelize probabilistic phenomena
    - Not direct relation with physical parameters



# Geometric models

- 1D Fire front on a 2D surface
- Huygens principle
- Advantages and disadvantages the same as  $\Rightarrow$  cellular automata.



# Empiric models

- An energy balance is considered on the fire front
- Rate of spread given by empiric laws
- R.C.Rothermel,  
*A Mathematical Model for Predicting Fire Spread in Wildland fuels.*
- Advantages  $\Rightarrow$  Fast computation
- Disadvantages  $\Rightarrow$  Parameters must be adjusted case by case.



# Combustion models: Convection-Difussion-Radiation-Reaction

## Complex models

- Several phases and conservation laws are considered
- Solid phase and gas phase with different temperature
- Two layer models
- Three dimensional equations
- Large time computation

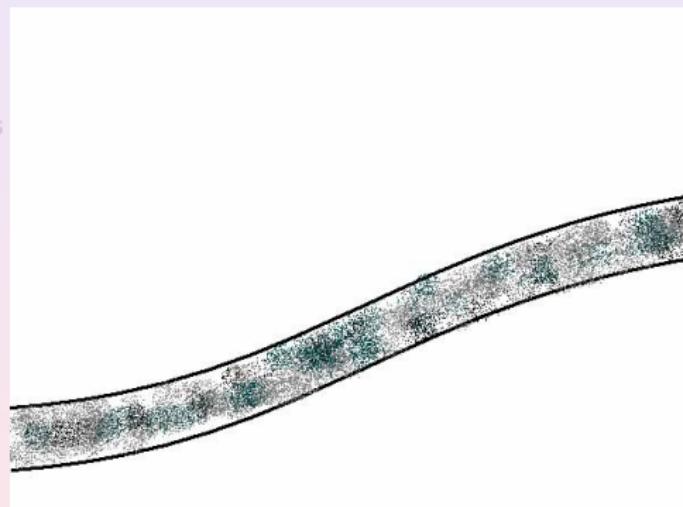
## Simplified models

- An average medium is considered
- Only one temperature
- One phase (other phases parametrized)
- One or two dimensional equations considered
- Could allow real time computation or faster



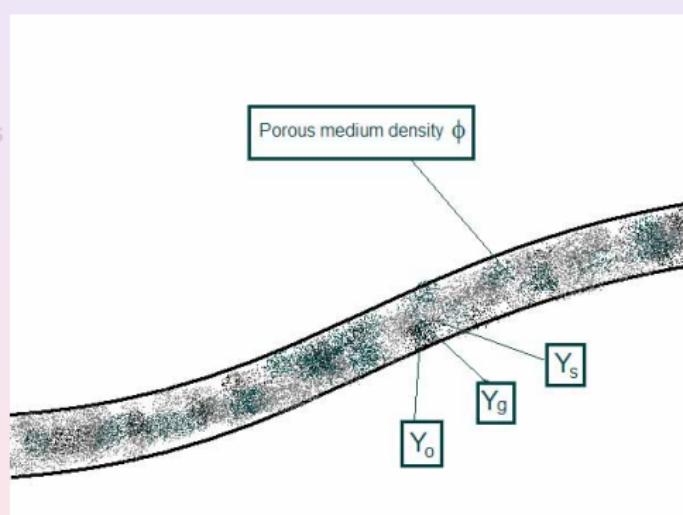
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- ② Composed of solid fuel  $Y_s$ , oxygen  $Y_o$  and gaseous fuel  $Y_g$
- ③ Gases are all assumed to have an equal temperature  $T_g$ . The solid temperature  $T_s$  is coupled by the term  $h(T_s - T_g)$
- ④ Solid fuel  $Y_s$  transforms into gaseous fuel  $Y_g$  through pyrolysis.
- ⑤ Gaseous fuel  $Y_g$  reacts with the oxygen  $Y_o$
- ⑥ Which generates the flames and heat
- ⑦ Gases and temperatures are under the influence of convection, diffusion and heat loss in the vertical direction (cooling)
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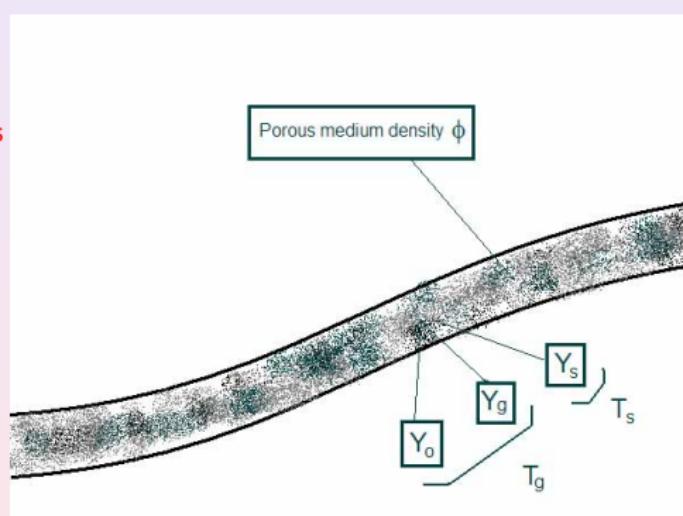
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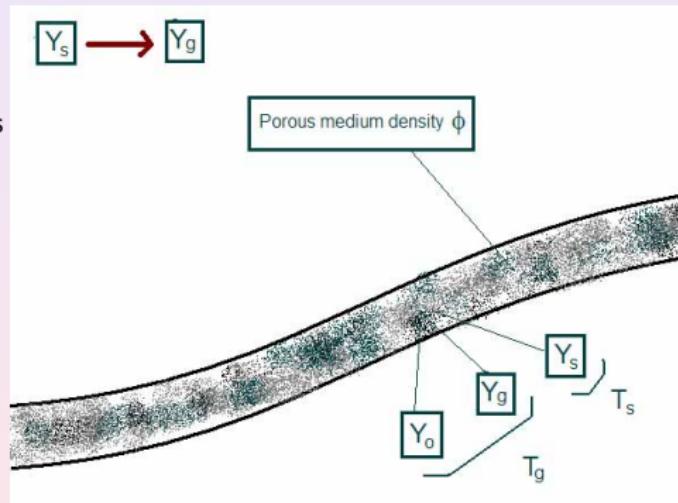
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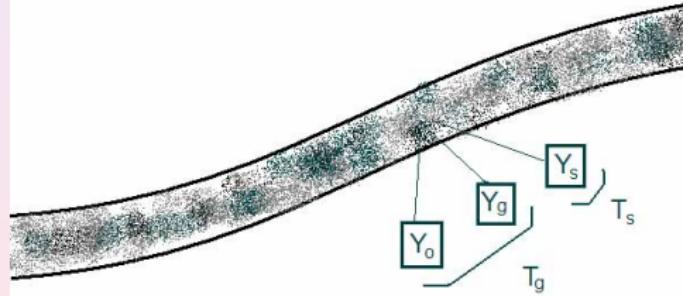
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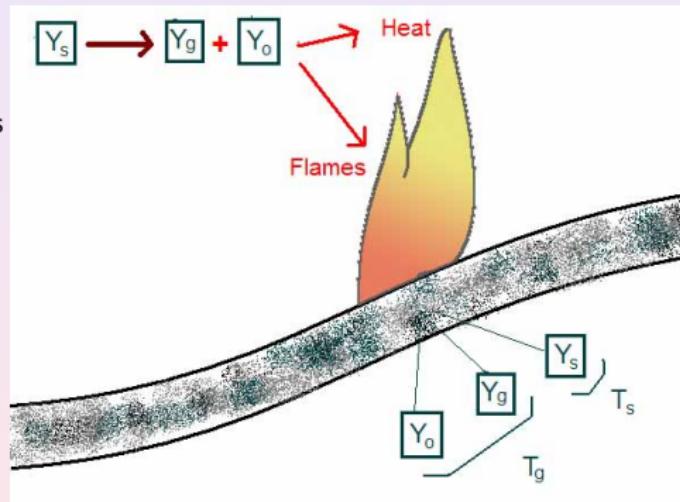
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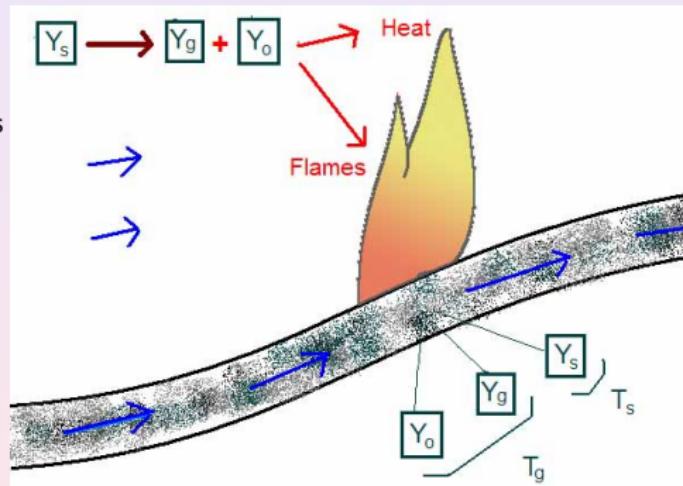
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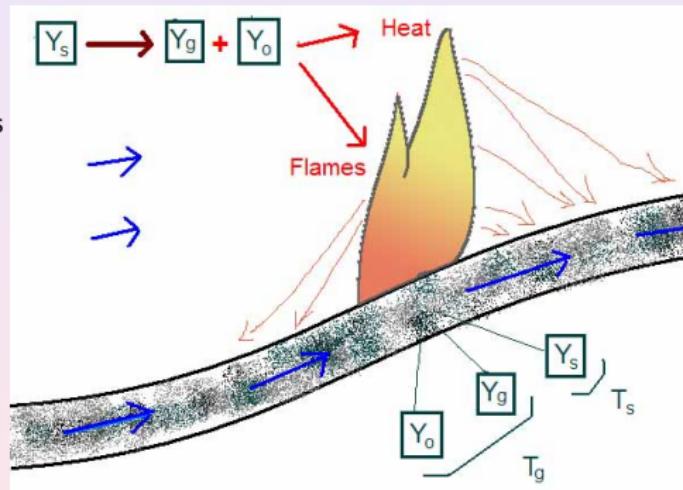
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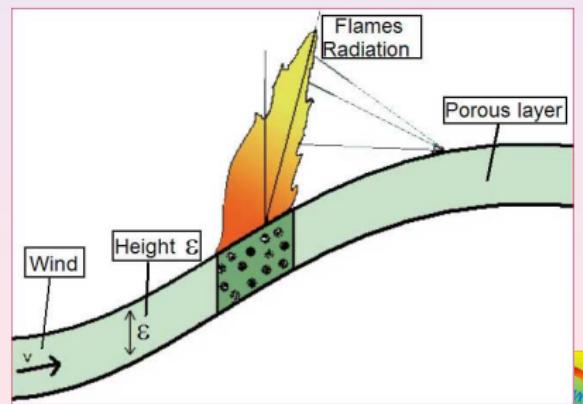
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- Convective transport due to the wind
- Diffusion within the vegetation
- Vertical cooling or vertical loss
- Link between  $T_s$  and  $T_g$
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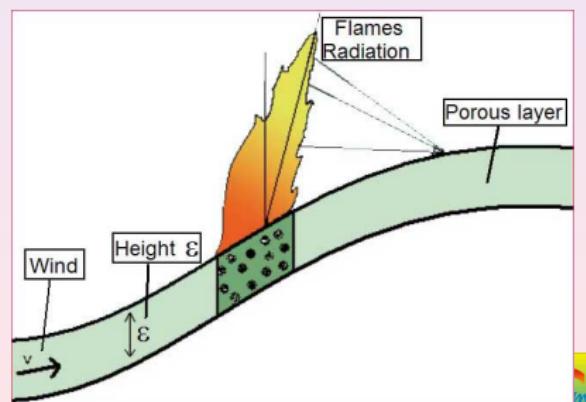
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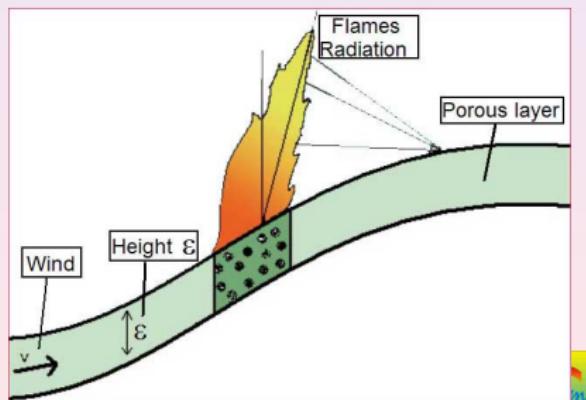
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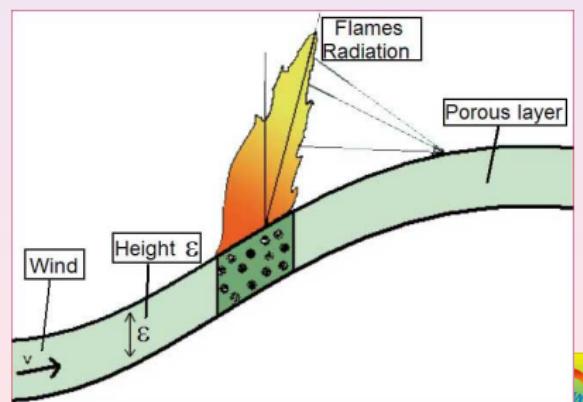
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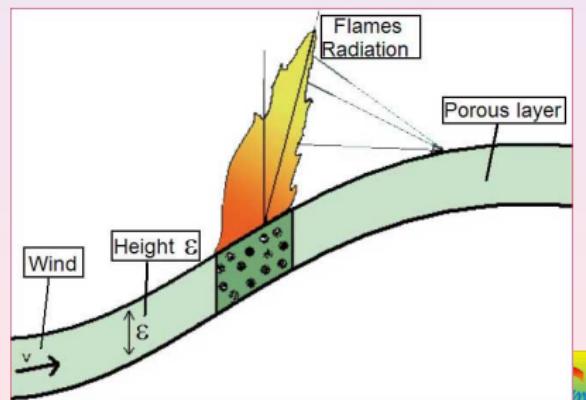
$$② \phi\rho C\left(\frac{\partial T_g}{\partial t} + \mathbf{v} \cdot \nabla T_g\right) - \nabla(\kappa_g \nabla T_g) = -h(T_g - T_s) + qA_g Y_o Y_g e^{-\frac{E_a}{RT_g}} - h_s(T_g - T_\infty)$$

$$③ \frac{\partial Y_s}{\partial t} = -A_s Y_s e^{-\frac{E_a}{RT_s}}$$

$$④ \quad \frac{\partial Y_o}{\partial t} + \mathbf{v} \cdot \nabla Y_o - \nabla(\kappa_{oo} \nabla Y_o) = -h_o(Y_o - Y_{o,\infty}) - A_g Y_o Y_g e^{-\frac{E_a}{RTg}}$$

$$⑤ \quad \frac{\partial Y_g}{\partial t} + \mathbf{v} \cdot \nabla Y_g - \nabla(\kappa_{gg} \nabla T_s) = -h_g(Y_g - 0) - A_g Y_o Y_g e^{-\frac{E_a}{RTg}} + A Y_s e^{-\frac{E_a}{RTs}}$$

- Time terms
  - Convective transport due to the wind
  - Diffusion within the vegetation
  - Vertical cooling or vertical loss
  - Link between  $T_s$  and  $T_g$
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  - Radiation term  $R(T_g)$



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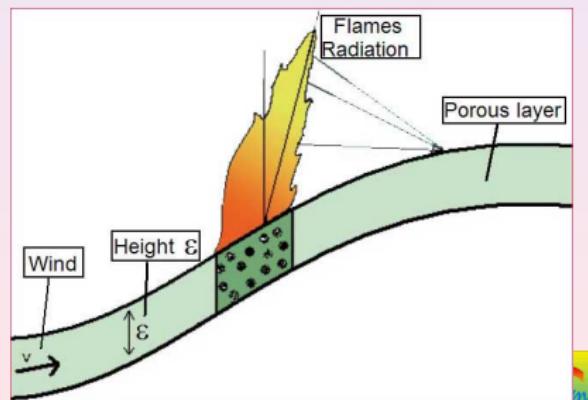
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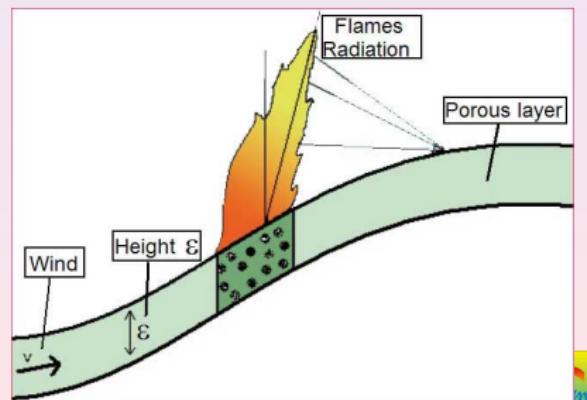
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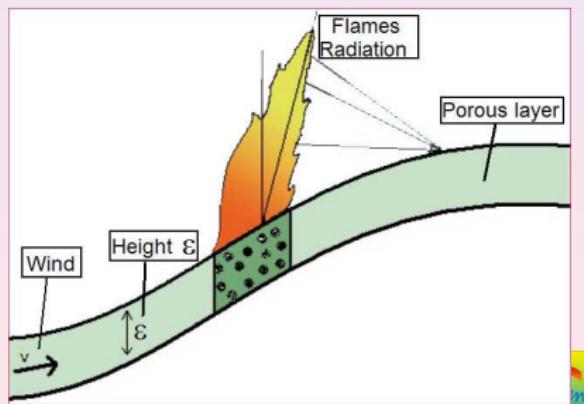
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- The goal
- A Simple Physical Model
- Enthalpy operator

## 3 NUMERICAL METHOD

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- Solution at each time step

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- Characteristic Method

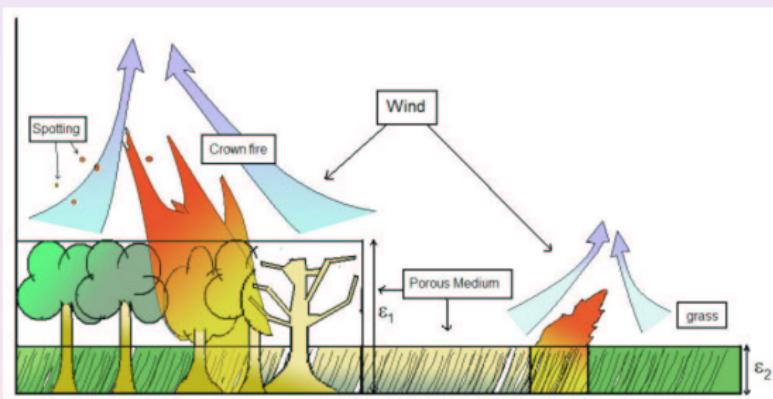
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- Physical and numerical data



Numerical Simulation of forest fires in small computers with computation times being a small fraction of the real time



# Requirements

## Simple models

- Simplified mathematical models
- Basically two dimensional models
- Some physical phenomena will be parametrized.

## Realistic models

- Take into account main mechanisms of propagation
- Take into account three dimensional effects



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# Convection-Diffusion-Radiation simplified models

Require the solution 1D or 2D of Convection-Diffusion-Radiation equations.

layer conducting radiating material

- $\rho C \left( \frac{\partial u}{\partial t} + \mathbf{V} \nabla u \right) - (J * u^4 - u^4) + \alpha u = Af(u, y)$
- $\frac{\partial y}{\partial t} = -f(u, y)$

- $u$  Temperature
- $y$  Mass fraction of fuel
- $-(J * u^4 - u^4)$  Non local diffusion (due to radiation)



# Local approximation

Estimation of the term  $-(J * u^4 - u^4)$

$$\begin{aligned}(J * u^4)(x) - u^4 &= \int J(x-y)u^4(y)dy - u^4(x) \\&= \int J(x-y)(u^4(y) - u^4(x))dy \\&= \int J(x-y)((u^4)'(x)(y-x) + \frac{1}{2}(u^4)''(x) + \dots)dy \\&= -(\int J(z)zdz)(u^4)'(x) + (\int J(z)z^2dz)(u^4)''(x) +\end{aligned}$$

If the kernel  $J(.)$  is symmetric then

$$-(J * u^4 - u^4) \approx -(\int J(z)z^2dz)(4u^3u')'(x)$$



If the kernel is not symmetric ( due to wind which tilt the flame)  
a convection term appears

$$-(J * u^4 - u^4) \approx \left( \int J(z)z dz \right) (u^4)'(x) - \left( \int J(z)z^2 dz \right) (4u^3 u')'(x)$$



# Diffusion approximation

If it is assumed that the absorption mean penetration distance  $\frac{1}{a}$  is small compared with the distances over which significant temperatures changes occur, that is the effect of radiation is local (Rosseland approximation)

layer conducting radiating material

- $\rho C \left( \frac{\partial u}{\partial t} + \mathbf{V} \nabla u \right) - \nabla \left( \frac{16\sigma u^3}{3a_R} \nabla u \right) + \alpha u = Af(u, y)$
- $\frac{\partial y}{\partial t} = -f(u, y)$

- $a_R$  Rosseland absorption

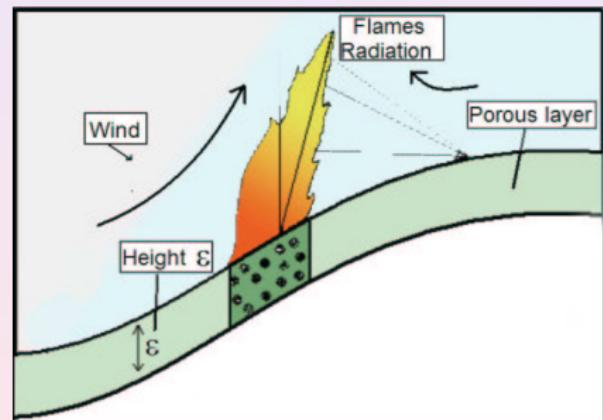


# A Simple Physical Model

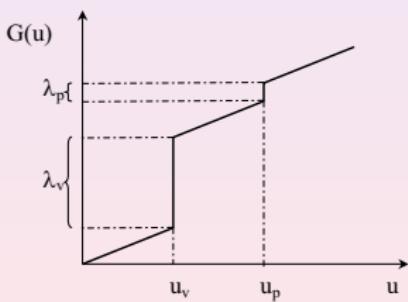
The non dimensional equations governing the fire spread in a region  $\Omega$  are,

- ①  $\partial_t e + \alpha u = r(u, y)$
- ②  $e \in G(u)$
- ③  $\partial_t y = -g(u)y$

- $r(u, y)$  is the source term
- $G(u)$  is the Enthalpy
- $g(u)y$  is the rate of pyrolysis



$$G(u) = \begin{cases} u & \text{if } u < u_v \\ [u_v, u_v + \lambda_v] & \text{if } u = u_v \\ u + \lambda_v & \text{if } u_v < u < u_p \\ [u_p + \lambda_v, u_p + \lambda_v + \lambda_p] & \text{if } u = u_p \\ u + \lambda_v + \lambda_p & \text{if } u > u_p \end{cases}$$



## Comments

- Local diffusion has been neglected
- Convection has been neglected but main **WIND EFFECTS** are considered via radiation model
- Main mechanism of fuel heating is **RADIATION** (non local diffusion)
- Effects of the **WATER CONTENT** are considered
- Gas phase is parametrized (The temperature and height of the flame are parameters)



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# Euler semi-implicit method

Let  $\Delta t = t^{n+1} - t^n$  a time step and let  $y^n$ ,  $e^n$  and  $u^n$  denote approximations at time step  $t^n$ , to the exact solution  $y$ ,  $e$  and  $u$  respectively.

We consider a semi-implicit scheme. At each time step we solve,

$$\frac{e^{n+1} - e^n}{\Delta t} + \alpha u^{n+1} = r^n$$
$$e^{n+1} \in G(u^{n+1})$$

$$\frac{y^{n+1} - y^n}{\Delta t} = -y^{n+1}g(u^{n+1})$$



# Solving the nonlinear equation in each time step

The multivalued operator  $G(.)$  is maximal monotone, then its resolvent  $J_\lambda = (Id + \lambda G)^{-1}$  for any  $\lambda > 0$  is a well defined univalued operator.

Taking  $\lambda = 1/(\alpha\Delta t)$  the algorithm simplifies

$$\textcircled{1} \quad u^{n+1} = J_{1/\alpha\Delta t} \left( \frac{1}{\alpha\Delta t} e^n + \frac{1}{\alpha} r^n \right)$$

$$\textcircled{2} \quad e^{n+1} = e^n - \alpha\Delta t u^{n+1} + \Delta t r^n$$

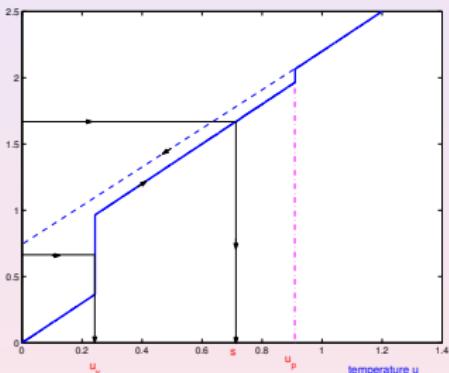
$$\textcircled{3} \quad y^{n+1} = \frac{y^n}{1 + \Delta t g(u^{n+1})}$$



# Practical Computation of the Resolvent

$$(\alpha\Delta t I + G)\bar{z} \ni \bar{u}$$

- if  $(1 + \alpha\Delta t)u_v < \bar{u} < (1 + \alpha\Delta t)u_v + \lambda_v$ ,



- if  $(1 + \alpha\Delta t)u_v + \lambda_v < \bar{u} < (1 + \alpha\Delta t)u_p + \lambda_v$ ,

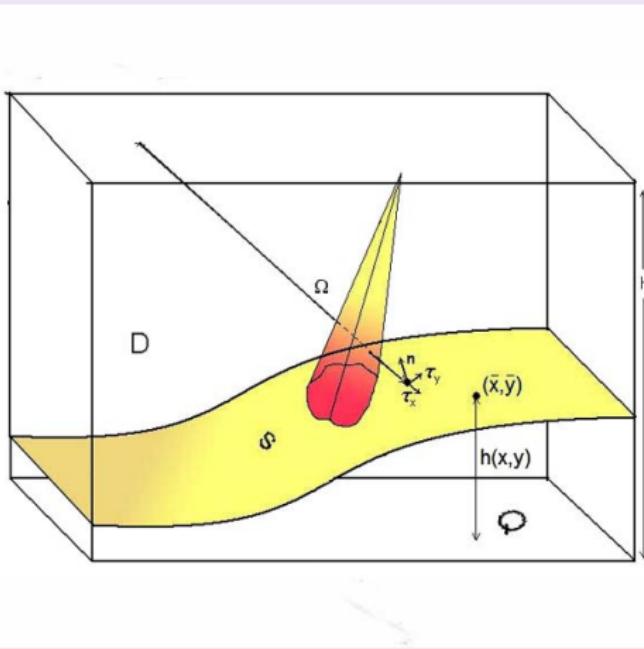
$$\bar{z} = \frac{\bar{u} - \lambda_v}{1 + \alpha\Delta t}$$

Figure: Computation of  $\bar{z}$



# Nonlocal radiation

The source term is due to nonlocal radiation: In each point it is calculated summing up the radiation intensity for all directions  $\Omega$ .



1- Assume as known  $i(\bar{x}, \Omega)$   $\forall \Omega$  for a given  $\bar{x}$

2- By use of polar coordinates at the tangent space of  $\bar{x}$  and defining  $\mu = \cos \theta$  and  $\gamma = \cos \phi$

$$r(\bar{x}) = \int i(\bar{x}, \Omega) \Omega \cdot n \, d\omega = \int_{\mu, \gamma} \frac{i_+(\bar{x}, \mu, \gamma) \mu}{\sqrt{1-\gamma^2}} \, d\mu d\gamma$$

3- A quadrature may be done with respect to:

$\mu \rightarrow$  Gauss-Legendre

$\gamma \rightarrow$  Gauss-Chevichev



# Radiation equation

After adimensionalization, the radiation equations in the direction  $\Omega$  can be written as

$$\begin{aligned}\Omega \cdot \nabla i + a^* i &= \delta(1 + u_g)^4 \quad \text{in } D \\ i &= 0 \quad \text{on } \partial D \cap \{\mathbf{x}; \Omega \cdot \mathbf{n} < 0\}\end{aligned}$$

The incident energy at a point  $\mathbf{x}(x, y, h(x, y))$  of the surface  $S$  per unit time and per unit area will be obtained summing up the contribution of all directions  $\Omega$

$$r(\mathbf{x}) = \int_{\omega=0}^{2\pi} i(\mathbf{x}, \Omega) \Omega \cdot \mathbf{n} \, d\omega \quad (1)$$



The radiation term  $r$  in the energy equation is computed by numerical integration.

Summing up for all the solid angles

$$r(\bar{x}) = \int_{\theta=0}^{\theta=\pi/2} \int_{\phi=0}^{\phi=2\pi} i(\bar{x}, \theta, \phi) \cos \theta \sin \theta \, d\theta d\phi = \\ \int_{\mu=0}^{\mu=1} \int_{\gamma=-1}^{\gamma=1} \frac{i_+(\bar{x}, \mu, \gamma)\mu}{\sqrt{1-\gamma^2}} \, d\mu d\gamma + \int_{\mu=0}^{\mu=1} \int_{\gamma=-1}^{\gamma=1} \frac{i_-(\bar{x}, \mu, \gamma)\mu}{\sqrt{1-\gamma^2}} \, d\mu d\gamma \quad (2)$$

$$\mu = \cos \theta, \gamma = \cos \phi$$

$i_+$  (resp.  $i_-$ ) stands for the radiation intensity  $i$  corresponding to an angle  $\phi$  such that  $0 \leq \phi < \pi$  (resp.  $\pi \leq \phi < 2\pi$ ).



The integrals are computed using Gauss-Legendre quadrature with respect  $\mu$  and Gauss-Chebyshev quadrature with respect  $\gamma$  in order to cope with the singular weight  $\frac{1}{\sqrt{1-\gamma^2}}$ .

$$r(\bar{x}) \approx \sum_{k,l} W_{kl} i_+(\bar{x}, \mu_k, \gamma_l) \mu_k + \sum_{k,l} W_{kl} i_-(\bar{x}, \mu_k, \gamma_l) \mu_k \quad (3)$$



# Numerical solution of the radiation equation

The characteristic line is

$$[0, \xi] \longrightarrow \mathcal{R}^3$$

$$\xi \longrightarrow (x(\xi) = \bar{x} + \xi\Omega_1, y(\xi) = \bar{y} + \xi\Omega_2, z(\xi) = \bar{z} + \xi\Omega_3)$$

On the characteristic, the radiation equation becomes

$$\frac{di}{d\xi} + a^*i = \delta(1 + u_g)^4$$

which can be solved backwards together with the condition

$$\lim_{\xi \rightarrow \infty} i(\xi) = 0 \quad (4)$$



# Extension of the temperature field

When integrating the radiation equation we need to know  $a^*(\xi(x, y, z))$  and  $u_g(\xi(x, y, z))$  so we need to extend the temperature field to the whole domain  $D$ .

- If there is **NO WIND**, we extend the temperature vertically

$$\tilde{u}(x, y, z) = u(x, y, h(x, y))$$

$$h(x, y) < z < h(x, y) + H$$

- In the case of **WIND CONDITIONS** we compute the extended field assuming a convective transport

$$\tilde{u}(x, y, z) = u\left(x - (z - h(x, y)) \frac{v_x}{v_z}, y - (z - h(x, y)) \frac{v_y}{v_z}, h(x, y)\right)$$



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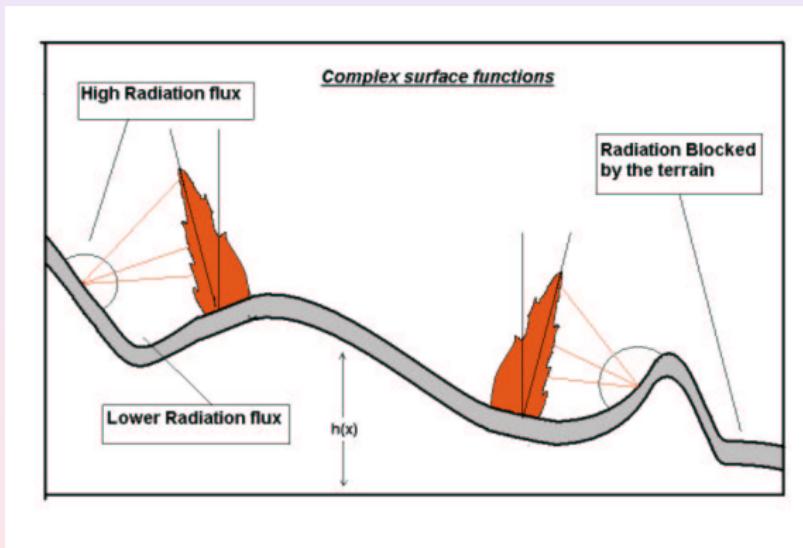
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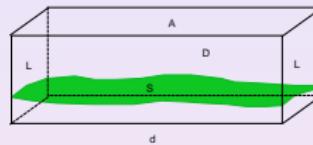


# The radiation equation and complex geometries

The method is able to cope with complex geometries



# Vertical diffusion model



Gives the horizontal wind field in a 3D domain by the expression

$$V(\mathbf{x}, z) = m(\mathbf{x}, z) \nabla p(\mathbf{x}) + n(\mathbf{x}, z) \nabla \hat{T}(\mathbf{x})$$

- ①  $m(\mathbf{x}, z) = \frac{1}{2}z^2 - \delta z - \frac{1}{2}h^2(\mathbf{x}) + (\delta + \xi)h(\mathbf{x}) - \xi\delta$
- ②  $n(\mathbf{x}, z) = -\frac{1}{24}z^4 + \frac{1}{6}\delta z^3 - \frac{1}{3}\delta^3 z + \frac{1}{24}h^4(\mathbf{x}) - \dots$
- ③  $p(\mathbf{x})$  is a potential function



# Vertical diffusion model

The potential  $p(\mathbf{x})$  satisfies the following boundary problem

$$\begin{aligned} -\nabla(a\nabla p) &= \nabla(b\nabla \hat{T}) \quad \text{in } d \\ a\frac{\partial p}{\partial n} &= -b\frac{\partial \hat{T}}{\partial \nu} + (\delta - h)v_m \cdot \nu \quad \text{on } \partial d \end{aligned}$$

- ①  $a(\mathbf{x}) = \frac{1}{3}(\delta - h(\mathbf{x}))^2(3\xi + \delta - h(\mathbf{x}))$
- ②  $b(\mathbf{x}) = \frac{1}{30}(\delta - h(\mathbf{x}))^2 \left( 2\delta^2(2\delta + 5\xi) - 2\delta(\delta - 5\xi)h(\mathbf{x}) - (3\delta + 5\xi)h^2(\mathbf{x}) + h^3(\mathbf{x}) \right)$



# Optimal control problem

The optimal control problem is characterised by



$$\int_{\omega} a \nabla p(u) \nabla \varphi + \frac{1}{\alpha} \int_{\partial\omega} q \varphi = - \int_{\omega} b \nabla \hat{T} \nabla \varphi \quad \forall \varphi \in V$$



$$\begin{aligned} & \int_{\omega} a \nabla q(u) \nabla \psi \\ & - \sum_{i=1}^N \int_{\omega} \rho_{\epsilon}(x - x_i) (m \nabla p(u) + n \nabla \hat{T} - V_i) m \nabla \psi = 0 \quad \forall \psi \in V \end{aligned}$$



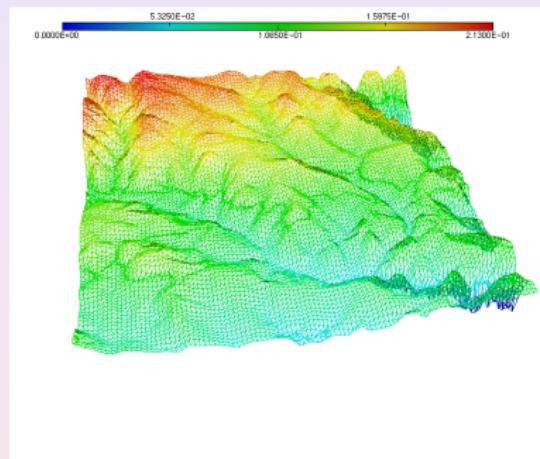
$$u = -\frac{1}{\alpha} q \quad \text{on } \partial d$$



# An example: Simulation: Cofrentes (Valencia, Spain)

## The data

- Zone 5 Km long by 5 Km width
- Average Meteorological wind given.
- The ignition point
- GIS data: Surface orography, vegetation density
- Combustion properties of the fuel: half-life decay.



Data by courtesy of  
Tecnosylva, S.L.  
León, Spain



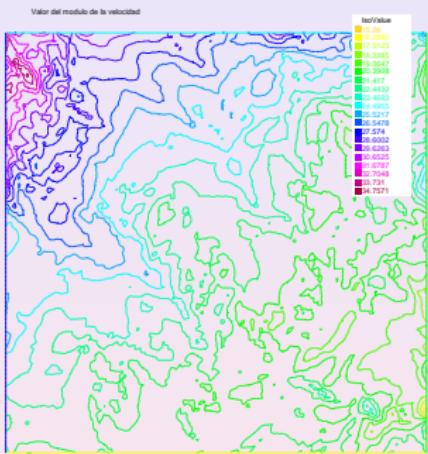


Figure: Wind Module

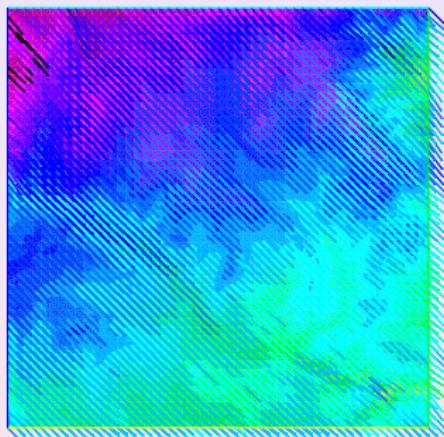


Figure: Wind direction

# Evolution of the fire front with Wind

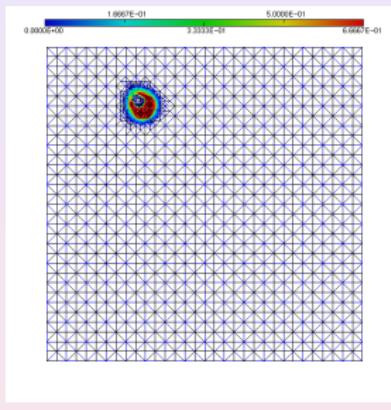


Figure: With wind at time 33 minutes

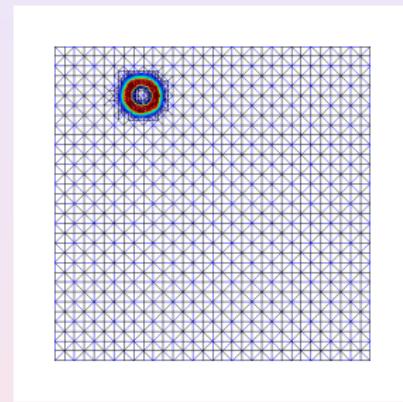


Figure: Without wind at time 33 minutes



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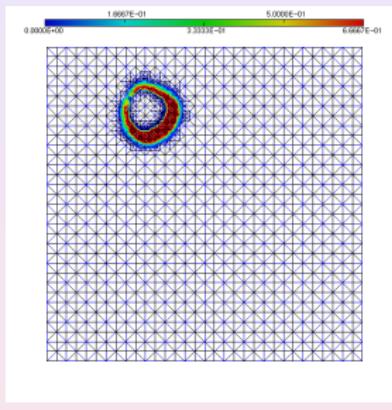


Figure: With wind at time 66 minutes

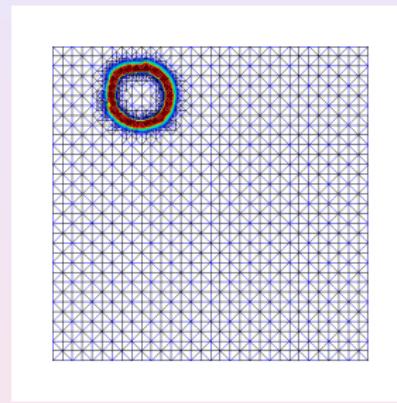


Figure: Without wind at time 66 minutes



# Evolution of the fire front with Wind

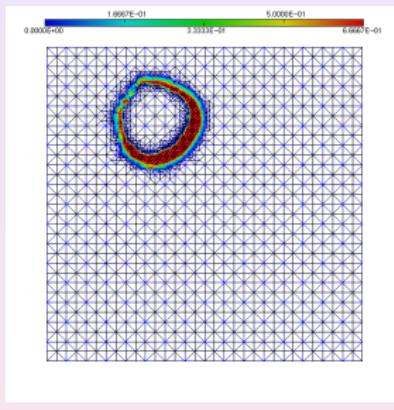


Figure: With wind at time 99 minutes

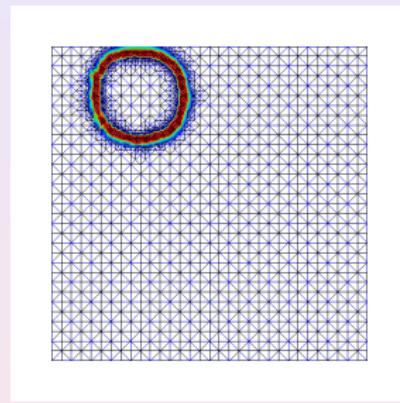


Figure: Without wind at time 99 minutes



# Evolution of the fire front with Wind

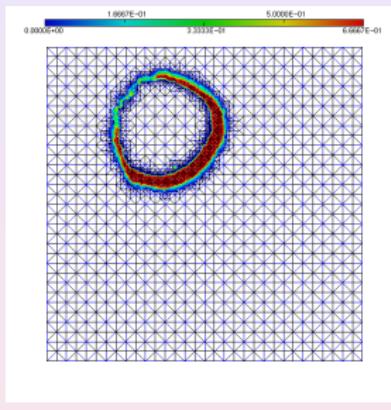


Figure: With wind at time 132 minutes

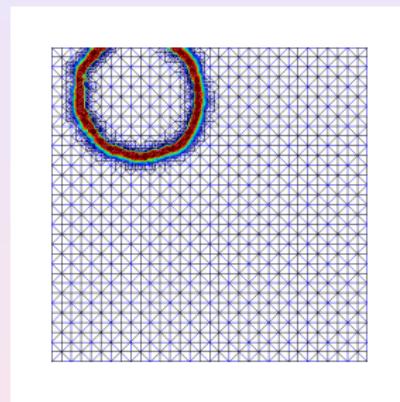


Figure: Without wind at time 132 minutes



# Evolution of the fire front with Wind

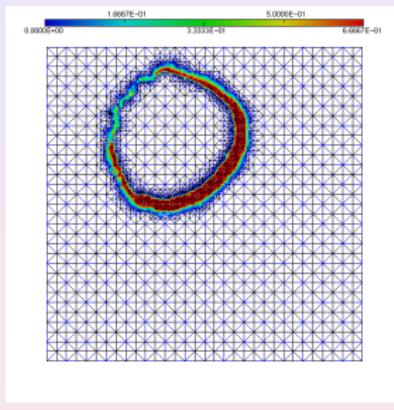


Figure: With wind at time 165 minutes

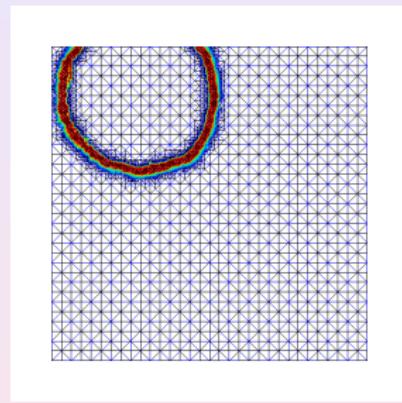


Figure: Without wind at time 165 minutes



# Evolution of the fire front with Wind

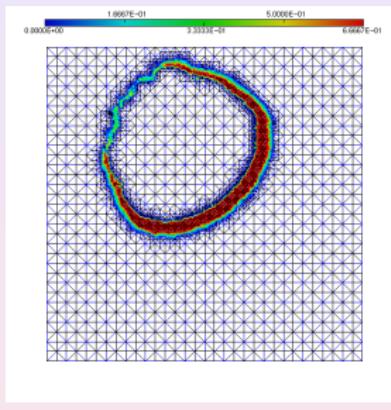


Figure: With wind at time 198 minutes

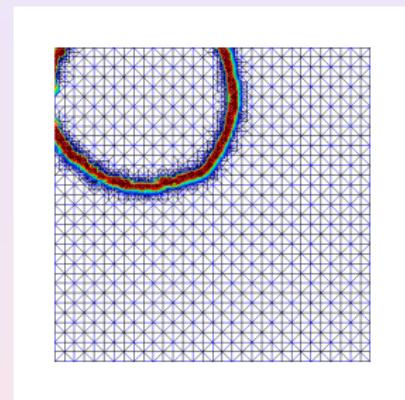


Figure: Without wind at time 198 minutes



# Evolution of the fire front with Wind

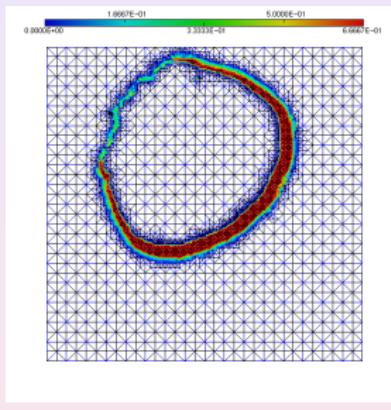


Figure: With wind at time 231 minutes

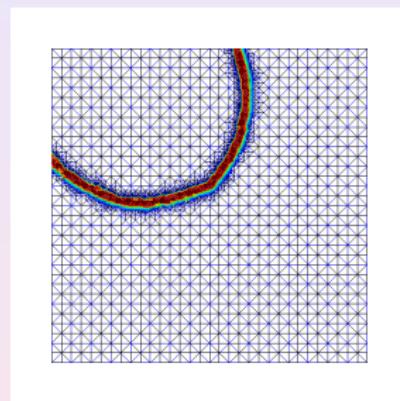


Figure: Without wind at time 231 minutes



# Evolution of the fire front with Wind

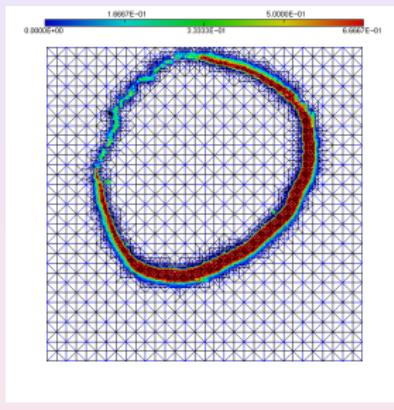


Figure: With wind at time 264 minutes

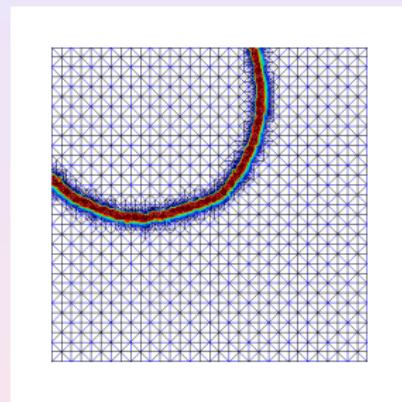


Figure: Without wind at time 264 minutes



# Evolution of the fire front with Wind

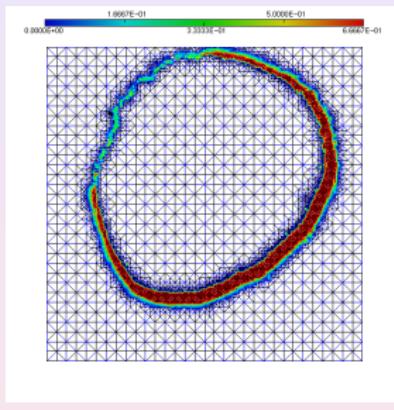


Figure: With wind at time 297 minutes

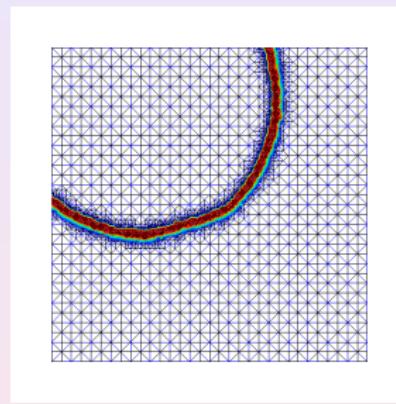


Figure: Without wind at time 297 minutes



# 3D representation



Figure: Without wind

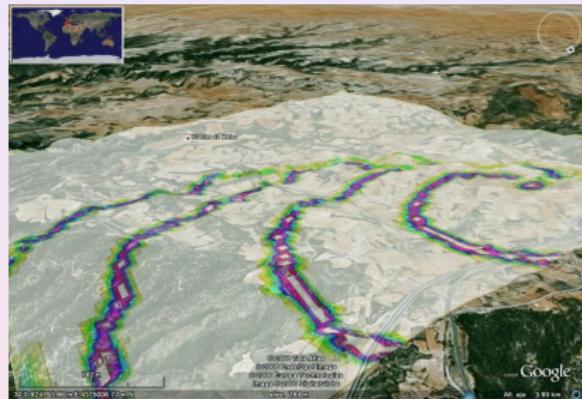


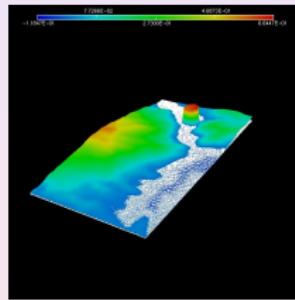
Figure: With wind



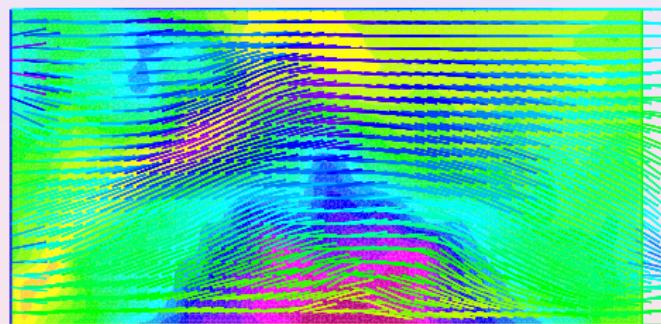
# An example: Simulation of a fire in a river basin

## The data

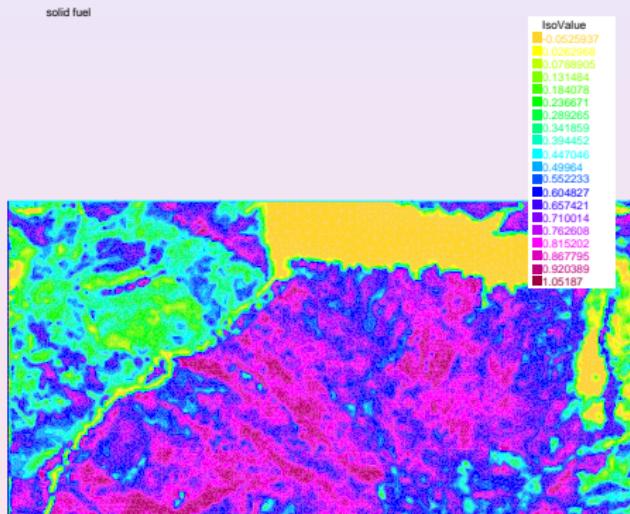
- Zone 6 Km long by 3 Km width
- Meteorological wind given in several points.
- The ignition point
- GIS data: Surface orography, vegetation density
- Combustion properties of the fuel: half-life decay.



# Wind: wind field $V$

velocidad,  $z=0.1+h$ 

# Initial Fuel Density



## Main physical parameters

- Zone 6 Km long by 3 Km width
- Water content: 2%
- Half-time decay: 700 seconds
- Flame height: 20 meters
- Flame temperature 1225C
- Wind velocity (1st): NO WIND
- Wind velocity (2nd):  $10m/s \div 20m/s$

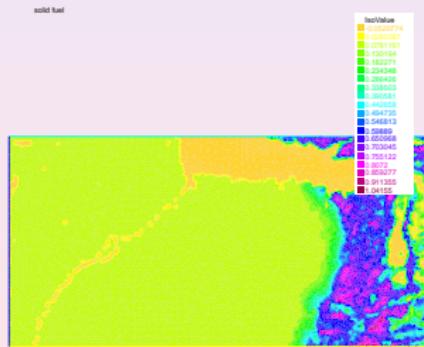
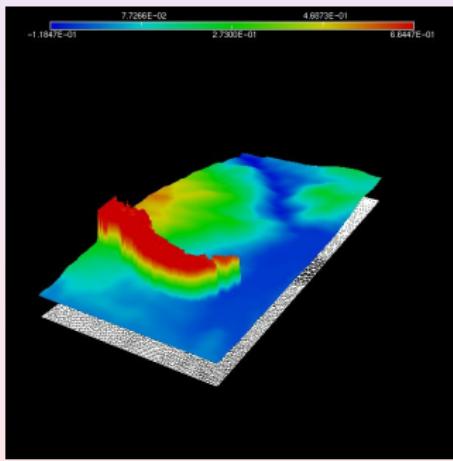
## Main numerical data

- Mesh size: maximum size 100 meters, minimum size 50 meters
- number of unknowns per variable: 31774
- Time step: 200 seconds
- Numerical integration points for radiation:  $8 \times 2 = 16$



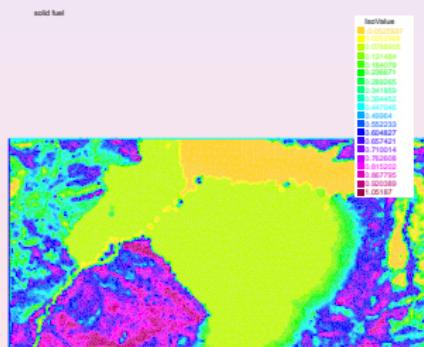
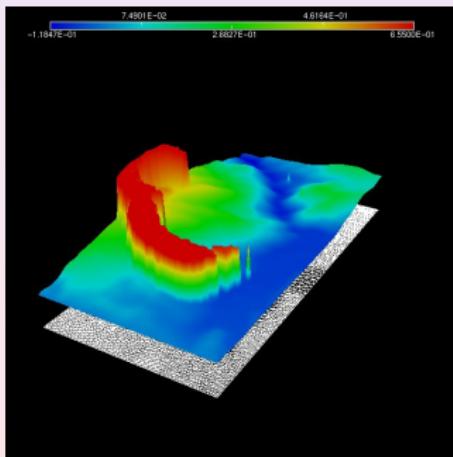
## Computation results

- Real time of the simulation  $\approx$  5 hours 30 minutes
- Computation time in my laptop: 25 minutes
- Mean fire front velocity: 650 m/h (**no wind**)



## Computation results

- Real time of the simulation  $\approx$  5 hours 30 minutes
- Computation time in my laptop: 25 minutes
- Mean fire front velocity: 710 m/h (**wind**)



# Evolution of a fire in the Ebro basin

SHOW ANIMATIONS

